Synaptic bistability due to nucleation and evaporation of receptor clusters

by

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V. M. Burlakov¹, N. Emptage², A. Goriely¹ and P. C. Bressloff³
¹Mathematical Institute, OCCAM, University of Oxford, 24-26 St Giles, Oxford OX1 3LB, UK
²Department of Pharmacology, University of Oxford, Mansfield Road, Oxford OX1 3QT, UK
³Department of Mathematics, University of Utah, Salt Lake City, UT 84112

We introduce a bistable mechanism for long-term synaptic plasticity based on switching between two metastable states that contain significantly different numbers of synaptic receptors. One state is characterized by a two-dimensional gas of mobile interacting receptors and is stabilized against clustering by a high nucleation barrier. The other state contains a receptor gas in equilibrium with a large cluster of immobile receptors, which is stabilized from growing further by the turnover rate of receptors into and out of the synapse. Transitions between the two states can be initiated by either an increase (potentiation) or a decrease (depotentiation) of the net receptor flux into the synapse. This changes the saturation level of the receptor gas and triggers nucleation or evaporation of receptor clusters.

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According to current understanding, the main biophysical mechanism for storing information in the central nervous system is activity-based changes in the strengths or efficacies of synaptic connections between neurons (synaptic plasticity) [1–4]. In many cases, synaptic plasticity is expressed by changes in the number of neurotransmitter protein receptors within the postsynaptic membrane of a stimulated neuron [5, 6]. However any modification in the number of independent receptors cannot fully account for long-term memory because of the limited dwell time of individual receptors, which are recycled into and out of a synapse several times an hour [7]. One suggested mechanism for stabilizing higher concentrations of receptors at a synapse is through interactions with submembrane scaffolding proteins [8–10]. An alternative mechanism has recently been proposed by Shouval [11], in which receptor clusters can survive much longer than individual receptors if the rate of receptor insertion into the membrane (exocytosis) is higher in the vicinity of other receptors due to receptor-receptor interactions and the rate of receptor removal from the membrane (endocytosis) is independent of such interactions. However, no particular mechanism for cluster formation has previously been proposed.

In this Letter we present a physical model of receptor stabilization that includes an explicit mechanism for receptor clustering. The latter is formulated along analogous lines to the well-known phenomenon of vapor-liquid phase transformations in supercritical liquid droplet nucleation and growth [12]. Vapour in our model is represented by mobile receptors within the postsynaptic membrane. We show that supercritical clusters can quickly be nucleated at a pre-existing nucleation site (heterogeneous nucleation [13]) if the surrounding vapor becomes dense enough, e.g. after an excitatory signal has stimulated the synapse. The nucleation site can be generated by any receptor that has been localized by interactions with sub-membrane scaffolding proteins. Once nucleated, the cluster grows until it reaches an equilibrium size determined by the balance between the rates of endocytosis and exocytosis. The cluster can also be evaporated if the receptor concentration in the vapor becomes too low, e.g. after an inhibitory signal. We show that such a mechanism of cluster formation/evaporation allows a synapse to exhibit bistability, i.e. there exist two distinct stable steady states with different efficacies (receptor numbers). Receptors in the state with low efficacy are predominantly in the vapor phase, while the efficacy of the other state is significantly increased by the formation of one or more receptor clusters. Transitions between the two stable states correspond to writing or erasing information. We also show that synaptic bistability is controlled by the total influx of receptors into the postsynaptic terminal and depends on the area of the postsynaptic membrane. Changing any of these parameters modifies the size and stability of the receptor cluster and modulates synaptic efficacy accordingly.

Model description. At most excitatory synapses in the brain, receptors are highly clustered at the postsynaptic density (PSD), which is the protein-rich domain in the postsynaptic membrane of a dendritic spine (see Fig. 1). The dendritic spine is a small (sub-micrometer) membranous extrusion that protrudes from a dendrite. Typically spines have a bulbous head that is connected to the parent dendrite through a thin spine neck. There are two main mechanisms for the transport of receptors between the PSD and extrasynaptic regions of the spine and dendrite [9]. First, surface receptors can undergo lateral membrane diffusion, in which the PSD acts as a spatially localized trap through receptor interactions with scaffolding proteins and the cytoskeleton. This is consistent with single-particle tracking experiments, which show surface receptors undergoing periods of free diffusion interspersed with periods of restricted motion in confinement domains that coincide with synapses [14–17]. Second, surface receptors may be internalized via endo-
cytosis and either stored within an intracellular pool or recycled to the surface via exocytosis [18].

In our model of receptor clustering we will focus on the second transport process and neglect lateral diffusion between the PSD and the extrasynaptic membrane (see Fig. 1). Thus, we assume that exocytosis and endocytosis occur directly in the PSD [19]. We also assume that surface receptors within the PSD are either mobile or part of a stationary cluster. At any time $t$ the total number of receptors of the PSD will depend on the balance between fluxes into and out of the postsynaptic membrane. The total receptor influx has two components $y = y_1 + y_2$, where $y_1$ is the influx due to local receptor recycling and $y_2$ is the contribution from an intracellular receptor pool triggered upon synaptic excitation (see Fig. 1). The outflow of receptors depends on the numbers of clustered and mobile receptors, with endocytosis rates determined by their local environments. One candidate mechanism for receptor clustering is short-range (Van der Waals) interactions between receptors with a characteristic energy of the order $k_B T$, where $T$ is room temperature [20]. This relatively low binding energy allows us to treat receptor clusters as two-dimensional liquid droplets, which are approximately circular as they minimize their surface energy. Hence, any receptor removal from the cluster does not change the cluster shape and integrity.

**Energy analysis.** Although receptor diffusion within the PSD is 10-50 times slower than within the extrasynaptic membrane [14-17], it is still fast enough to establish thermodynamic equilibrium of the receptor vapor. This allows us to define the chemical potential of receptors in the cluster $\mu_{cl}$ and in the vapor $\mu_{vap}$ as

$$\mu_{cl} = -\varepsilon_{cl} - \varepsilon_{adh} + e^{a/R} \gamma,$$
$$\mu_{vap} = k_B T \ln(x_{vap} \Lambda^2) - \varepsilon_{adh}, \quad (1)$$

where $\varepsilon_{cl}$ is the binding energy of receptors in the cluster, $R$ is the cluster radius, $\varepsilon_{adh}$ is the adhesive energy of receptors in the membrane, $\gamma$ is the cluster surface energy per receptor, $a$ is the spacing between receptors in the cluster, and $\Lambda$ is the receptor thermal wavelength in the vapor [21]. If we explicitly modeled receptor diffusion using a lattice gas model (see Ref. [22]) then the thermal wave length would be equal to $a$.

The vapor concentration $x_{vap}$ can be obtained by equating the receptor influx $y$ with the receptor outflow at steady state, and assuming that the area $S_0$ occupied by the mobile receptors (vapor) is approximately independent of the size of the cluster. The outflow of receptors contains two contributions: one from the clustered receptors and the other from the mobile ones. The balance of receptor fluxes in and out of the PSD is

$$y = \pi R^2 \nu e^{-\varepsilon_{cl}} e^{-\varepsilon_{adh} + \gamma/R} + S_0 x_{vap} \nu e^{-\varepsilon_{adh}}, \quad (2)$$

where, for convenience, we express all lengths in $a$-units and all energies in $k_B T$-units. The product $P_{vap} = \nu e^{-\varepsilon_{adh}}$ is the endocytosis rate of receptors in the vapor, with $\nu$ the attempt frequency (taken to be the same for both mobile and bound receptors). Similarly, the product $P_{cl} = \nu e^{-\varepsilon_{adh} - \varepsilon_{cl} + \gamma/R}$ is the endocytosis rate from the cluster, which assumes that the interaction between receptors naturally increases the dwell time for bound synaptic receptors (see also [14]) in accordance with their binding energy $\varepsilon_{cl}$. Note that this is entirely opposite to the approach used in Ref. [11], where receptor removal rates were assumed to be the same for mobile and bound receptors. Rearranging Eq. (2) we obtain the vapor concentration as a function of cluster size:

$$x_{vap}(R) = \frac{y}{S_0 P_{vap}} - \frac{\pi R^2}{S_0} e^{-\varepsilon_{cl} + \gamma/R}. \quad (3)$$

In the limit $R \rightarrow 1$ (in $a$-units), we can neglect the second term on the RHS of Eq. (3) so that $x_{vap}(0) = y/(S_0 P_{vap})$. Using this expression and writing the clustering radius $R = R(n_{cl}) \equiv \sqrt{n_{cl}/\pi}$ in terms of receptor number in the cluster $n_{cl}$ we obtain the energy required for formation of the receptor cluster from the vapor:

$$\Delta E = \int_0^{n_{cl}} (\mu_{cl} - \mu_{vap}) dn + S_0[F(n_{cl}) - F(0)]$$
$$= -n_{cl} \varepsilon_{cl} - \int_0^{n_{cl}} \ln \left[ x_{vap}(n_{cl}) - n_{cl} e^{-\varepsilon_{cl} + \gamma \sqrt{\pi n_{cl}}} \right] dn$$
$$+ 2 \gamma \sqrt{\pi n_{cl}} + S_0[F(n_{cl}) - F(0)], \quad (4)$$

where $F(n) = x_{vap}(R(n)) \ln(x_{vap}(R(n)))$. Fig. 2 shows the dependence of the system energy upon the cluster size.
tosis rate $P_{vap}$ and synaptic area $S_0$. The value of $\gamma$ is chosen in order to describe correctly the surface energy of relatively small discrete cluster of receptors on a hexagonal grid (see section 1 in Supporting Material). Also note that the receptor influx is increased up to $y = 0.6$; 2) $P_{vap} = 0.1, S_0 = 500, y = 0.9$; 3) $P_{vap} = 0.1, S_0 = 500, y = 0.5$; 4) $P_{vap} = 0.12, S_0 = 500, y = 0.6$; 5) $P_{vap} = 0.1, S_0 = 300, y = 0.36$. The latter value of $y$ is chosen to keep $x_{vap}(0)$ the same as in 1.

size for different values of the receptor influx $y$, endocytosis rate $P_{vap}$ and synaptic area $S_0$. The value of $\gamma$ is chosen in order to describe correctly the surface energy of a state with a cluster of size $n_{cl}$ and mobile $N_{vap}$ receptors:

$$\frac{dn_{cl}}{dt} = \frac{2\sqrt{\pi n_{cl}} (J_{con} - J_{evap})}{n_{cl}P_{cl}}$$

$$\frac{dN_{vap}}{dt} = y - N_{vap} P_{vap} - 2\sqrt{\pi n_{cl}} (J_{con} - J_{evap}) - n_{cl} P_{cl}$$

where $J_{con}$ and $J_{evap}$ are, respectively, the fluxes of condensing and evaporating receptors from the cluster. We take $J_{con} = \beta x_{vap} D / 4$, where $D/(4\lambda)$ is the rate at which receptors collide with the cluster surface, and $\beta$ is the probability for these receptors to be incorporated into the cluster. Here $D$ is the diffusion coefficient of mobile receptors in the PSD, and $\lambda$ is the receptor mean free path. For hopping-like diffusion in a lattice-gas model, $\lambda$ can be taken equal to $a$. The factor $1/4$ takes into account the fraction of receptors moving towards the cluster surface. We have made an implicit assumption that the receptor vapor is in thermodynamic equilibrium, i.e. homogeneous, which is well justified if $\beta \ll 1$. The flux $J_{evap} = \beta x_{GT} D / 4$ can be obtained by assuming that the cluster evaporates with the same rate as if it was in equilibrium with the surrounding vapor, and its stability is due to the compensating condensation flux from the vapor. The so-called Gibbs-Thomson concentration $x_{GT}$ of the vapor can be found from the equilibrium condition $\mu_{vap} = \mu_{con}$ with the chemical potentials given by Eqs (1). Hence, $x_{GT} = e^{-\varepsilon_{cl} + \gamma \sqrt{\pi / n_{cl}}}$.

We solved Eqs. (5) using kinetic Monte Carlo simulations (see section 2 of Supporting Material and Ref. [25]). Typical simulation results illustrating responses to an increased receptor influx (potentiation) or receptor removal rate (depotentiation) are presented in Fig. 3. In the simulations we used $a = 5 nm$, $\beta = 0.005$ and a characteristic diffusion coefficient $D$ for mobile receptors in the PSD [14–17]. We also assumed that the PSD contains a single nucleation site with at least two receptors. The impact of the parameter $\beta D$ on the cluster size and its fluctuations is described in section 3 of the Supporting Material. We find that our simulation results are in excellent agreement with the energy analysis if $\beta D > 80 a^2 / s$. The values of other parameters used for the simulations shown in Fig. 3a are the same as those used to calculate curve 1 in Fig. 2. Cluster formation was triggered by a short increase in $y_2$ from 0 to 1.2, after which the cluster grows to its stable size of around 200 receptors, and fluctuates about this stable value. To initiate cluster evaporation we increased $P_{vap}$ by factor of 3. The results shown in Fig. 3b were obtained using parameter values for curve 5 in Fig. 2. Increasing $y_2$
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References:
[19] There is strong experimental evidence that insertion and removal of surface receptors occurs at extrasynaptic sites so that lateral diffusion is required for receptors to enter or exit the PSD [18]. For simplicity, we lump these processes together. However, it would be straightforward to extend the model to include the effects of surface diffusion between the PSD and extrasynaptic sites.
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